# SPECTRUM OF A $k$ th-ORDER SLANT HANKEL OPERATOR 

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#### Abstract

For an integer $k \geq 2$, a $k$ th-order slant Hankel operator $D_{\phi}$ with symbol $\phi$ in $L^{\infty}$ is an operator on $L^{2}$ whose representing matrix is $\left(a_{-k i-j}\right)_{i, j=-\infty}^{\infty}$, where $a_{t}$ 's are the Fourier coefficients of $\phi$. In this paper, we discuss some spectral properties of the operator $D_{\phi}$ and identify its spectrum and essential spectrum.


## 1. Introduction

Ever since the introduction of the class of Hankel operators in the middle of twentieth century, it has assumed tremendous importance due to its far reaching applications to problems of rational approximation, information and control theory, interpolation and prediction problems etc. Hankel operators and their matrix variants have occured in realization problem for certain discrete time linear systems and in determining which systems are exactly controllable [10].

Various generalizations of the notion of Hankel operators have been the subject of research of many mathematicians. R.A.M. Avendano [4], [6] initiated the notion of $\lambda$-Hankel operators in the year 2000. Later in the year 2002, motivated by the work of Barria and Halmos, Avendano [5] brought the concept of essentially Hankel operators into picture. S.C. Arora [1], along with his research associates, discussed a new variant of the class of Hankel operators and introduced the class of slant Hankel operators. Subsequently, the notion of slant Hankel operators was generalized to $k$ th-order slant Hankel operators ( $k \geq 2, k$ an integer) by Arora and Bhola [2, [3].

In this paper, we carry this study further and investigate some spectral properties of $k$ th-order slant Hankel operators on the space $L^{2}(\mathbb{T}), \mathbb{T}$ denoting the unit circle. We begin with the following:

Let $\phi(z)=\sum_{i=-\infty}^{\infty} a_{i} z^{i}$ be a bounded measurable function on the unit circle $\mathbb{T}$, where $a_{i}=\left\langle\phi, z^{i}\right\rangle$ is the $i$ th Fourier coefficient of $\phi$ and $\left\{z^{i}: i \in \mathbb{Z}\right\}$ is the standard basis of the space $L^{2}(\mathbb{T})$.

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For any integer $k \geq 2$, let $W_{k}$ be an operator on $L^{2}$ such that

$$
W_{k}\left(z^{i}\right)= \begin{cases}z^{i / k}, & \text { if } i \text { is divisible by } k \\ 0 & \text { otherwise }\end{cases}
$$

In particular for $k=2$, the operator $W_{2}$ denoted by $W$ is

$$
\begin{aligned}
& W\left(z^{2 i}\right)=z^{i} \\
& W\left(z^{2 i-1}\right)=0
\end{aligned}
$$

for all $i$ in $\mathbb{Z}$. Let $J$ denote the reflection operator on $L^{2}$, that is $J: L^{2} \rightarrow L^{2}$, $J(f(z))=f(\bar{z})$. It is easy to see the following:
(i) $\left\|W_{k}\right\|=1,\|J\|=1$
(ii) $J W_{k}=W_{k} J$.
(iii) $J=J^{*}, J^{2}=I$
(iv) $W_{k}^{*} z^{i}=z^{k i}$, for all $i$ in $\mathbb{Z}$.
(v) $W_{k} W_{k}^{*}=I$ and $W_{k}^{*} W_{k}=P_{k}$, where $P_{k}$ denotes the projection of $L^{2}$ onto the closed linear span of $\left\{z^{k i}: i \in \mathbb{Z}\right\}$

Definition 1 ([1]). A slant Hankel operator $K_{\phi}$ induced by $\phi$ in $L^{\infty}$ on the space $L^{2}$ is defined as

$$
K_{\phi}=J W M_{\phi}
$$

where $J$ and $W=W_{2}$ are as defined above and $M_{\phi}$ denotes the multiplication operator induced by $\phi$ in $L^{\infty}$ on the space $L^{2}$.

It is known that [1] an operator $A$ on $L^{2}$ is a slant Hankel operator if and only if $M_{\bar{z}} A=A M_{z^{2}}$.

Arora and Bhola [2] generalized the notion of slant Hankel operators to $k$ th-order slant Hankel operators as follows:

Definition $2(\boxed{2})$. For an integer $k \geq 2$, a $k$ th-order slant Hankel operator $D_{\phi}$ on the space $L^{2}$ is defined as

$$
D_{\phi}=J W_{k} M_{\phi}
$$

where $J$ and $W_{k}$ are as defined above and $M_{\phi}$ is the multiplication operator on $L^{2}$ induced by $\phi$ in $L^{\infty}$. It is clear that a slant Hankel operator is a particular case of a $k$ th-order slant Hankel operator (for $k=2$ ).

For $\phi(z)=\sum_{i=-\infty}^{\infty} a_{i} z^{i} \in L^{\infty}$, the matrix representation of $D_{\phi}$ with respect to the standard basis $\left\{z^{i}: i \in \mathbb{Z}\right\}$ of $L^{2}$ is given by:

So,

$$
\begin{aligned}
\left\langle D_{\phi} z^{j}, z^{i}\right\rangle & =\left\langle J W_{k} M_{\phi} z^{j}, z^{i}\right\rangle \\
& =\left\langle W_{k} M_{\phi} z^{j}, z^{-i}\right\rangle \\
& =\left\langle\phi, z^{-k i-j}\right\rangle \\
& =a_{-k i-j}
\end{aligned}
$$

Thus, $D_{\phi}: L^{2} \rightarrow L^{2}$ is defined as

$$
D_{\phi}\left(z^{j}\right)=\sum_{t=-\infty}^{\infty} a_{-k t-j} z^{t}, \text { for all } j \text { in } \mathbb{Z}
$$

For a fixed integer $k \geq 2$, the set of all $k$ th-order slant Hankel operators is denoted by $\mathcal{D}_{k}[2]$. It is proved that [2] a bounded linear operator $D$ on the space $L^{2}$ is a $k$ th-order slant Hankel operator if and only if $M_{\bar{z}} D=D M_{z^{k}}$. Also it is shown that the set $\mathcal{D}_{k}$ is neither an algebra, nor a self adjoint set and contains no non-zero hyponormal operator, no non-zero compact operator and no isometry.In continuation we also observe here that $D_{\phi}$ in $\mathcal{D}_{k}$ is a coisometry if and only if

$$
\left|\phi\left(\frac{\theta}{k}\right)\right|^{2}+\left|\phi\left(\frac{\theta+2 \pi}{k}\right)\right|^{2}+\ldots+\left|\phi\left(\frac{\theta+(k-1) 2 \pi}{k}\right)\right|^{2}=k .
$$

For if $f$ is in $L^{2}$, then

$$
\begin{aligned}
\left\|D_{\phi}^{*} f\right\|_{2}^{2}= & \left\|M_{\bar{\phi}} W_{k}^{*} J f\right\|_{2}^{2} \\
= & \int_{0}^{2 \pi}|\phi(\theta)|^{2}|f(-k \theta)|^{2} \frac{d \theta}{2 \pi} \\
= & \int_{0}^{2 k \pi}\left|\phi\left(\frac{\theta}{k}\right)\right|^{2}|f(-\theta)|^{2} \frac{d \theta}{2 k \pi} \\
= & \frac{1}{k} \int_{0}^{2 \pi}\left\{\left|\phi\left(\frac{\theta}{k}\right)\right|^{2}+\left|\phi\left(\frac{\theta+2 \pi}{k}\right)\right|^{2}+\ldots\right. \\
& \left.+\left|\phi\left(\frac{\theta+(k-1) 2 \pi}{k}\right)\right|^{2}\right\}|f(-\theta)|^{2} \frac{d \theta}{2 \pi} \\
= & \left\|M_{\psi} \hat{f}\right\|_{2}^{2}
\end{aligned}
$$

where

$$
\psi=\left[\frac{1}{k}\left\{\left|\phi\left(\frac{\theta}{k}\right)\right|^{2}+\left|\phi\left(\frac{\theta+2 \pi}{k}\right)\right|^{2}+\ldots+\left|\phi\left(\frac{\theta+(k-1) 2 \pi}{k}\right)\right|^{2}\right\}\right]^{\frac{1}{2}}
$$

and $\hat{f}=J f$. Since for any $f$ in $L^{2},\left\|M_{\psi} f\right\|_{2}=\|f\|_{2}$ if and only if $|\psi|=1$ a.e on $\mathbb{T}$, it follows that $D_{\phi}$ is a coisometry iff

$$
\left|\phi\left(\frac{\theta}{k}\right)\right|^{2}+\left|\phi\left(\frac{\theta+2 \pi}{k}\right)\right|^{2}+\ldots+\left|\phi\left(\frac{\theta+(k-1) 2 \pi}{k}\right)\right|^{2}=k
$$

So far, only the algebraic properties and some topological properties of the set $\mathcal{D}_{k}$ have been studied. In this paper we intend to discuss the spectral properties of the operator $D_{\phi}$ and determine the spectrum and the essential spectrum of the operator $D_{\phi}$ for some special symbols $\phi$ in $L^{\infty}$. To determine the desired ones we firstly make the following observations:

Observation 1: Any operator unitarily equivalent to a slant Toeplitz operator is also a slant Toeplitz operator.

If $T$ is a slant Toeplitz operator, then $M_{z} T=T M_{z^{2}}$. Now, for any unitary operator $U$ on $L^{2}$, we have

$$
\begin{aligned}
\left(U^{*} M_{z} U\right)\left(U^{*} T U\right) & =U^{*} M_{z} T U \\
& =U^{*} T M_{z^{2}} U \\
& =\left(U^{*} T U\right)\left(U^{*} M_{z} U\right)^{2}
\end{aligned}
$$

Since any operator unitarily equivalent to a shift operator is again a shift operator of the same multiplicity, it follows that $U^{*} T U$ is also a slant Toeplitz operator. Observation 2: A non-zero slant Toeplitz operator cannot be a slant Hankel operator.

For if, $A$ is a slant Toeplitz operator on $L^{2}$ which is also slant Hankel, then the matrix $\left(\alpha_{i j}\right)$ of $A$ with respect to the standard orthonormal basis satisfies

$$
\begin{aligned}
& \alpha_{i j}=\alpha_{i+1, j+2} \\
& \alpha_{i j}=\alpha_{i-1, j+2}
\end{aligned}
$$

for all $i, j \in \mathbb{Z}$. As $A$ is a bounded operator on $L^{2}$, we have $\alpha_{i j}=0$ for all $i, j \in \mathbb{Z}$.
From the above two observations, it follows that a non-zero slant Toeplitz operator cannot be unitarily equivalent to a slant Hankel operator.

## 2. Spectrum of $D_{\phi}$

It is interesting to note that though there is no unitary equivalence between a slant Toeplitz operator and a kth-order slant Hankel operator, yet some of the spectral properties are the same as is shown in this section. In fact here, we aim at proving that if $\phi$ in $L^{\infty}$ is invertible then the spectrum of $D_{\phi}$ contains a closed disc. More specifically, if $\phi$ is an inner function then the spectrum of $D_{\phi}$ contains the closed unit disc. To accomplish the task we begin with the following:

Lemma 3. For an invertible $\phi$ in $L^{\infty}, \sigma_{p}\left(D_{\phi}\right)=\sigma_{p}\left(D_{\phi\left(z^{k}\right)}\right)$, where $\sigma_{p}\left(D_{\phi}\right)$ denotes the point spectrum of $D_{\phi}$.

Proof. Suppose that $\lambda \in \sigma_{p}\left(D_{\phi}\right)$. Then there exists a non-zero $f$ in $L^{2}$ such that $D_{\phi} f=\lambda f$. Let $F=J(\phi f)$. Then $F$ is in $L^{2}$ and

$$
\begin{aligned}
D_{\phi\left(z^{k}\right)} F & =D_{\phi\left(z^{k}\right)}(J \phi f) \\
& =J W_{k}\left(\phi\left(z^{k}\right) J \phi f\right) \\
& =J \phi(z) \cdot J W_{k}(\phi f) \\
& =J \phi(z) \cdot D_{\phi} f \\
& =J \phi(z) \cdot \lambda f \\
& =\lambda J \phi f \\
& =\lambda F .
\end{aligned}
$$

Therefore, $\lambda \in \sigma_{p}\left(D_{\phi\left(z^{k}\right)}\right)$. Conversely, let $\mu \in \sigma_{p}\left(D_{\phi\left(z^{k}\right)}\right)$. Then there exists a non-zero $g$ in $L^{2}$ satisfying $D_{\phi\left(z^{k}\right)} g=\mu g$. Let $G=\phi^{-1} J g$. Then $G$ is in $L^{2}$ and $G$
satisfies

$$
\begin{aligned}
D_{\phi} G & =J W_{k} M_{\phi}\left(\phi^{-1} J g\right) \\
& =J W_{k} J g \\
& =\phi^{-1} \phi W_{k} g \\
& =\phi^{-1} W_{k}\left(\phi\left(z^{k}\right) g\right) \\
& =\phi^{-1} J D_{\phi\left(z^{k}\right)} g \\
& =\phi^{-1} J \mu g \\
& =\mu G .
\end{aligned}
$$

Thus, $\mu$ is in $\sigma_{p}\left(D_{\phi}\right)$. This proves the lemma.
Lemma 4. For any $\phi$ in $L^{\infty}, \sigma\left(D_{\phi}\right)=\sigma\left(D_{\phi\left(z^{k}\right)}\right)$, where $\sigma\left(D_{\phi}\right)$ denotes the spectrum of $D_{\phi}$.
Proof. It is known that [7 for any two bounded linear operators $A$ and $B$ on a Hilbert space $H$,

$$
\sigma(A B) \cup\{0\}=\sigma(B A) \cup\{0\}
$$

It follows that

$$
\begin{aligned}
\sigma\left(D_{\phi}^{*}\right) \cup\{0\} & =\sigma\left(M_{\bar{\phi}} W_{k}^{*} J\right) \cup\{0\} \\
& =\sigma\left(W_{k}^{*} J M_{\bar{\phi}}\right) \cup\{0\} \\
& =\sigma\left(J W_{k}^{*} M_{\bar{\phi}}\right) \cup\{0\} \\
& =\sigma\left(J^{2} U_{\phi\left(z^{k}\right)}^{*} J\right) \cup\{0\} \\
& =\sigma\left(D_{\phi\left(z^{k}\right)}^{*}\right) \cup\{0\} .
\end{aligned}
$$

Here, $U_{\phi}$ denotes the operator $W_{k} M_{\phi}$ on the space $L^{2}$. From this we get,

$$
\sigma\left(D_{\phi}\right) \cup\{0\}=\sigma \overline{\left(D_{\phi}^{*}\right)} \cup\{0\}=\sigma \overline{\left(D_{\phi\left(z^{k}\right)}^{*}\right)} \cup\{0\}=\sigma\left(D_{\phi\left(z^{k}\right)}\right) \cup\{0\}
$$

Also we see that the operator $J W_{k}^{*}$ is not onto as the range of $J W_{k}^{*}$ is the closed linear span of $\left\{\bar{z}^{k n}: n \in \mathbb{Z}\right\}$ in $L^{2}$. Thus,

$$
c l .\left\{\text { Range }\left(J W_{k}^{*} M_{\bar{\phi}}\right)\right\} \neq L^{2}
$$

Equivalently, there is a non-zero $f$ in $\operatorname{Ker}\left(W_{k}^{*} M_{\bar{\phi}}\right)$. That is, $M_{\bar{\phi}\left(z^{k}\right)} W_{k}^{*} J(J f)=0$, where $0 \neq J f$ is in $L^{2}$. Thus, $0 \in \sigma_{p}\left(D_{\phi\left(z^{k}\right)}^{*}\right)$ which in turn implies $0 \in \sigma_{p}\left(D_{\phi\left(z^{k}\right)}\right)$. Now, if $\phi$ is invertible, then by Lemma $3,0 \in \sigma_{p}\left(D_{\phi}\right)$ and the conclusion is immediate. If $\phi$ is not invertible then we can have a sequence $\left\{\phi_{n}\right\}$ of invertible functions in $L^{\infty}$ such that $\left\|\phi_{n}-\phi\right\|_{\infty} \rightarrow 0$ as $n \rightarrow \infty$, for the set $\left\{\phi \in L^{\infty}: \phi^{-1} \in L^{\infty}\right\}$ is dense in $L^{\infty}$ [8]. Since each $\phi_{n}$ is invertible, therefore $0 \in \sigma_{p}\left(D_{\phi_{n}}\right)$ for every $n$. Hence there exists $f_{n} \neq 0$, corresponding to each $n$, such that $D_{\phi_{n}} f_{n}=0$. Without loss of generality we assume that $\left\|f_{n}\right\|=1$ and we have

$$
\begin{aligned}
\left\|D_{\phi} f_{n}\right\| & =\left\|D_{\phi} f_{n}-D_{\phi_{n}} f_{n}+D_{\phi_{n}} f_{n}\right\| \\
& \leq\left\|D_{\phi} f_{n}-D_{\phi_{n}} f_{n}\right\|+\left\|D_{\phi_{n}} f_{n}\right\| \\
& \leq\left\|\phi-\phi_{n}\right\|_{\infty} \\
& \rightarrow 0
\end{aligned}
$$

as $n \rightarrow \infty$. Therefore, $0 \in \pi\left(D_{\phi}\right)$, the approximate point spectrum of $D_{\phi}$ and hence $0 \in \sigma\left(D_{\phi}\right)$. This completes the proof.

Theorem 5. Let $D_{\phi} \in \mathcal{D}_{k}$, where $\phi$ in $L^{\infty}$ is invertible. Then $\sigma\left(D_{\phi}\right)$ contains a closed disc. Furthermore, the radius of this closed disc is $\frac{1}{r\left(D_{\psi}\right)}$, where $\psi=(J \bar{\phi})^{-1}$ and $r\left(D_{\psi}\right)$ denotes the spectral radius of $D_{\psi}$.
Proof. Let $\lambda \neq 0$ and assume that $D_{\phi\left(z^{k}\right)}^{*}-\lambda I$ is onto. For any $f$ in $L^{2}$ we have:

$$
\begin{aligned}
\left(D_{\phi\left(z^{k}\right)}^{*}-\lambda\right) f & =\left(J W_{k} M_{\phi\left(z^{k}\right)}\right)^{*} f-\lambda f \\
& =\bar{\phi}\left(z^{k}\right) \cdot W_{k}^{*} J f-\lambda f \\
& =W_{k}^{*}(\bar{\phi} J f)-\lambda f \\
& =W_{k}^{*}(\bar{\phi} J f)-\lambda\left(P_{k} f \oplus P_{k}^{\prime} f\right) \\
& =\left(W_{k}^{*}(\bar{\phi} J f)-\lambda P_{k} f\right) \oplus\left(-\lambda P_{k}^{\prime} f\right) \\
& =\lambda W_{k}^{*} M_{\bar{\phi}}\left(\lambda^{-1}-M_{\bar{\phi}^{-1}} W_{k}\right)(J f) \oplus\left(-\lambda P_{k}^{\prime} f\right) \\
& =\lambda W_{k}^{*} M_{\bar{\phi}}\left(\lambda^{-1}-J\left(W_{k} J M_{\bar{\phi}^{-1}\left(z^{k}\right)}\right)(J f) \oplus\left(-\lambda P_{k}^{\prime} f\right)\right) \\
& =\lambda W_{k}^{*} M_{\bar{\phi}}\left(\lambda^{-1}-J W_{k} M_{\bar{\phi}^{-1}\left(z^{k}\right)}\right)(J f) \oplus\left(-\lambda P_{k}^{\prime} f\right) \\
& =\lambda W_{k}^{*} M_{\bar{\phi}}\left(\lambda^{-1}-D_{\bar{\phi}^{-1}\left(\bar{z}^{k}\right)}\right)(J f) \oplus\left(-\lambda P_{k}^{\prime} f\right)
\end{aligned}
$$

where $P_{k}$ denotes the projection of $L^{2}$ on the closed linear span of $\left\{z^{k n}: n \in \mathbb{Z}\right\}$ and $P_{k}^{\prime}=I-P_{k}$.

Now let $0 \neq g_{0} \in P_{k}^{\prime}\left(L^{2}\right)$. Since $D_{\phi\left(z^{k}\right)}^{*}-\lambda I$ is onto, therefore there exists a non-zero $f$ in $L^{2}$ such that

$$
\left(D_{\phi\left(z^{k}\right)}^{*}-\lambda I\right) f=g_{0}
$$

This implies that

$$
\lambda W_{k}^{*} M_{\bar{\phi}}\left(\lambda^{-1}-D_{\bar{\phi}^{-1}\left(\bar{z}^{k}\right)}\right)(J f) \oplus\left(-\lambda P_{k}^{\prime} f\right)=g_{0}
$$

Since $g_{0} \in P_{k}^{\prime}\left(L^{2}\right)$, we have

$$
\lambda W_{k}^{*} M_{\bar{\phi}}\left(\lambda^{-1}-D_{\bar{\phi}^{-1}\left(\bar{z}^{k}\right)}\right)(J f)=0 .
$$

As $\lambda \neq 0, W_{k}^{*}$ is an isometry and $M_{\bar{\phi}}$ is invertible, therefore we have

$$
\begin{aligned}
& \left(\lambda^{-1}-D_{\bar{\phi}^{-1}\left(\bar{z}^{k}\right)}\right)(J f)=0 \\
& \Rightarrow \lambda^{-1} \in \sigma_{p}\left(D_{\bar{\phi}^{-1}\left(\bar{z}^{k}\right)}\right) .
\end{aligned}
$$

Since for any $\lambda \in \rho\left(D_{\phi\left(z^{k}\right)}^{*}\right)$, the resolvent set of $D_{\phi\left(z^{k}\right)}^{*}$, we have $D_{\phi\left(z^{k}\right)}^{*}-\lambda I$ is invertible and hence onto, above calculations yield

$$
\begin{aligned}
\mathfrak{D}=\left\{\lambda^{-1}: \lambda \in \rho\left(D_{\phi\left(z^{k}\right)}^{*}\right)\right\} & \subseteq \sigma_{p}\left(D_{\bar{\phi}^{-1}\left(\bar{z}^{k}\right)}\right) \\
& =\sigma_{p}\left(D_{(J \bar{\phi})^{-1}\left(z^{k}\right)}\right) .
\end{aligned}
$$

Replacing $(J \bar{\phi})^{-1}$ by $\phi$ we get that for any invertible $\phi$ in $L^{\infty}$, the spectrum of $D_{\phi}$ contains a disc of eigenvalues of $D_{\phi}$. Since the spectrum of any bounded linear
operator is compact, $\sigma\left(D_{\phi}\right)$ contains a closed disc. Also,

$$
\begin{aligned}
\max \left\{\left|\lambda^{-1}\right|: \lambda \in \rho\left(D_{\phi\left(z^{k}\right)}^{*}\right)\right\} & =\left[\min \left\{|\lambda|: \lambda \in \rho\left(D_{\phi\left(z^{k}\right)}^{*}\right)\right\}\right]^{-1} \\
& =\left[r\left(D_{\phi\left(z^{k}\right)}^{*}\right)\right]^{-1} \\
& =\left(r\left(D_{\phi\left(z^{k}\right)}\right)\right)^{-1} \\
& =\left[r\left(D_{\phi}\right)\right]^{-1}
\end{aligned}
$$

Replacing $\phi$ by $(J \bar{\phi})^{-1}$, we have that the radius of the closed disc contained in $\sigma\left(D_{\phi}\right)$ is $\frac{1}{r\left(D_{\psi}\right)}$, where $\psi=(J \bar{\phi})^{-1}$.

From the proof of the above theorem, it follows that

$$
\left[r\left(D_{\psi}\right)\right]^{-1} \leq r\left(D_{\phi}\right), \text { where } \psi=(J \bar{\phi})^{-1}
$$

Now if $\phi$ in $L^{\infty}$ is unimodular then

$$
\begin{aligned}
\left\|D_{\phi}^{n}\right\|^{2} & =\left\|D_{\phi}^{n} D_{\phi}^{* n}\right\| \\
& =\|I\| \\
& =1
\end{aligned}
$$

so that $r\left(D_{\phi}\right)=\lim _{n \rightarrow \infty}\left\|D_{\phi}^{n}\right\|^{1 / n}=1$. Moreover, if $|\phi|=1$ then $\left|\bar{\phi}^{-1}\right|=1$. Therefore in this case

$$
\left[r\left(D_{\psi}\right)\right]^{-1} \leq 1 \text { and } 1 \leq\left[r\left(D_{\psi}\right)\right]^{-1}
$$

Therefore $\sigma\left(D_{\phi}\right)$ contains the closed unit disc.
Hence we have the following:
Theorem 6. If $|\phi|=1$, then $\sigma\left(D_{\phi}\right)=\overline{\mathfrak{D}}$ where $\overline{\mathfrak{D}}$ denotes the closed unit disc in the complex plane. In particular if $\phi$ is an inner function then $\sigma\left(D_{\phi}\right)$ contains $\overline{\mathfrak{D}}$.

Remark 7. From the proof of Theorem 5, it may be seen that if $\phi$ is invertible in $L^{\infty}$ then $\sigma\left(D_{\phi}\right)$ contains a disc of eigenvalues of infinite multiplicity. Therefore in this case $\sigma_{e}\left(D_{\phi}\right)$, the essential spectrum of $D_{\phi}$ also contains a closed disc. So we have the following:
(i) For any invertible $\phi$ in $L^{\infty}$,

$$
\mathfrak{D} \subset \sigma_{e}\left(D_{\phi}\right) \subset \sigma\left(D_{\phi}\right)
$$

(ii) If $\phi$ is unimodular, then

$$
\sigma_{e}\left(D_{\phi}\right)=\sigma\left(D_{\phi}\right)=\overline{\mathfrak{D}}
$$

where $\overline{\mathfrak{D}}$ denotes the closed unit disc in the complex plane.
Theorem 8. Let $D_{\phi} \in \mathcal{D}_{k}$, where $\phi$ in $L^{\infty}$ is unimodular. If $|\lambda|<1$, then $\lambda$ is an eigenvalue of $D_{\phi}$ such that the corresponding eigenspace is generated by

$$
\sum_{n=0}^{\infty} \lambda^{n} D_{\phi}^{*^{n}} f
$$

where $f \in \operatorname{span}\left\{\bar{\phi} \bar{z}^{n k-j}: n \in \mathbb{Z}, j=1,2, \ldots, k-1\right\}$

Proof. Let $\phi \in L^{\infty}$ be unimodular. Then $D_{\phi}^{*}$ is an isometry. Therefore, by von Neumann-Wold decomposition for isometries [9], we can write

$$
D_{\phi}^{*} \simeq U \oplus V
$$

where $U$ is unitary on $\bigcap_{n=0}^{\infty} D_{\phi}^{*^{n}}\left(L^{2}\right)$ and $V$ is unilateral shift on $\bigoplus_{n=0}^{\infty} D_{\phi}^{*^{n}}\left(L^{2}\right)$.
Since $D_{\phi} \simeq U^{*} \oplus V^{*}$ and the eigenvalues of $V^{*}$ are $\{\lambda:|\lambda|<1\}$, it follows that $\{\lambda:|\lambda|<1\}$ is a set of eigenvalues of $D_{\phi}$ also.

For $|\lambda|<1$, consider the function $g=\sum_{n=0}^{\infty} \lambda^{n} D_{\phi}^{*^{n}} f$, where $0 \neq f \in\left(D_{\phi}^{*}\left(L^{2}\right)\right)^{\perp}$. We have,

$$
\begin{aligned}
|g(z)| & =\left|\sum_{n=0}^{\infty} \lambda^{n} D_{\phi}^{*^{n}} f\right| \\
& \leq \sum_{n=0}^{\infty}\left|\lambda^{n}\right|\left\|D_{\phi}^{*^{n}} f\right\| \\
& <\infty
\end{aligned}
$$

for $|\lambda|<1$ and $\left\|D_{\phi}^{*^{n}} f\right\|=\|f\|<\infty$.
Also,

$$
\begin{aligned}
D_{\phi} g & =D_{\phi} f+\lambda \sum_{n=1}^{\infty} \lambda^{n-1} D_{\phi}^{*^{n-1}} f \\
& =0+\lambda g \\
& =\lambda g
\end{aligned}
$$

Therefore, $g$ is an eigenvector corresponding to the eigenvalue $\lambda$ of $D_{\phi}$. Further,

$$
\begin{aligned}
D_{\phi}^{*}\left(L^{2}\right) & =\left(J W_{k} M_{\phi}\right)^{*}\left(L^{2}\right) \\
& =M_{\bar{\phi}} W_{k}^{*} J\left(L^{2}\right) \\
& =\operatorname{span}\left\{\bar{\phi} \bar{z}^{n k}: n \in \mathbb{Z}\right\}
\end{aligned}
$$

If follows that $\left(D_{\phi}^{*}\left(L^{2}\right)\right)^{\perp} \supseteq$ span $\left\{\bar{\phi} \bar{z}^{n k-j}: n \in \mathbb{Z}, j=1,2, \ldots, k-1\right\}$ because for any $m, n$ in $\mathbb{Z}$ and $j=1,2, \ldots, k-1$,

$$
\begin{aligned}
\left\langle\bar{\phi} \bar{z}^{n k}, \bar{\phi} \bar{z}^{m k-j}\right\rangle & =\left\langle\phi \bar{\phi}, \bar{z}^{m k-n k-j}\right\rangle \\
& =\left\langle 1, \bar{z}^{(m-n) k-j}\right\rangle \\
& =0
\end{aligned}
$$

for $(m-n) k-j \neq 0$ for any $m$ and $n$ in $\mathbb{Z}, j=1,2, \ldots, k-1$. On the other hand, since $\bar{\phi}$ is invertible, $\bar{\phi} L^{2}=L^{2}$. Therefore,

$$
\operatorname{span}\left\{\bar{\phi} \bar{z}^{n k}: n \in \mathbb{Z}\right\} \oplus \operatorname{span}\left\{\bar{\phi} \bar{z}^{n k-j}: n \in \mathbb{Z}, j=1,2, \ldots, k-1\right\}=\bar{\phi} L^{2}=L^{2}
$$

so that

$$
\left(D_{\phi}^{*}\left(L^{2}\right)\right)^{\perp} \subseteq \operatorname{span}\left\{\bar{\phi} \bar{z}^{n k-j}: n \in \mathbb{Z}, j=1,2, \ldots k-1\right\}
$$

Hence,

$$
\left(D_{\phi}^{*}\left(L^{2}\right)\right)^{\perp}=\operatorname{span}\left\{\bar{\phi} \bar{z}^{n k-j}: n \in \mathbb{Z}, j=1,2, \ldots k-1\right\}
$$

It follows that the eigenspace corresponding to the eigenvalue $\lambda$ (where $|\lambda|<1$ ) of $D_{\phi}$ is generated by

$$
\sum_{n=0}^{\infty} \lambda^{n} D_{\phi}^{*^{n}} f
$$

where $0 \neq f \in \operatorname{span}\left\{\bar{\phi} \bar{z}^{n k-j}: n \in \mathbb{Z}, j=1,2, \ldots k-1\right\}$.
This completes the proof.
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